



## Chap 9 Stability in the Frequency Domain

林沛群  
國立台灣大學  
機械工程學系

### Introduction

#### ▫ Roodmap

- ◆ Chapter 6 : Routh-Hurwitz Criterion

Check stability of the system by examining  $\Delta(s)$

Introduce idea of relation stability

- ◆ Chapter 7 : Root Locus

Investigate loci of the system poles as the system parameter changes

- ◆ Chapter 8 : Polar plot & Bode plot

Introduce frequency response of the system

- ◆ This Chapter :

Investigate stability of the system in the frequency domain

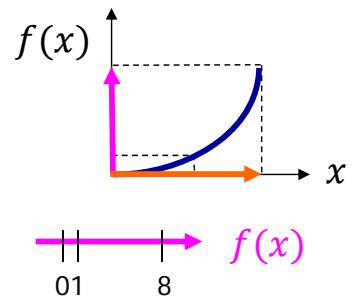
## Mapping Contours in the s-plane

### □ Real-number functions

$$x, f(x) \in \mathbb{R}$$

$$\text{Ex: } f(x) = x^3$$

$$f(x) : x \rightarrow f(x)$$



### □ Contour map

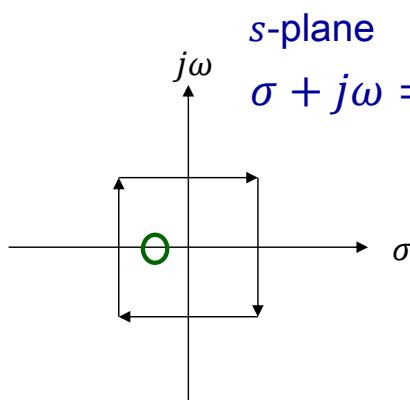
- ◆ A contour or trajectory in one plane mapped or translated into another plane by a relation  $F(s)$

$$\begin{array}{ccc} F(s) : s \rightarrow F(s) & & \\ \downarrow & & \downarrow \\ \sigma + j\omega & u + jv & \end{array}$$

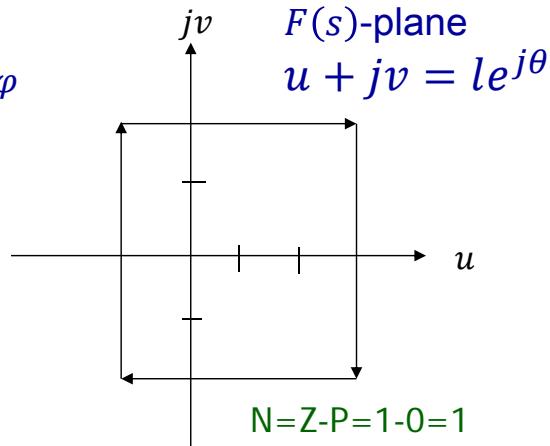
## Mapping Contours in the s-plane

### □ Example: $F(s) = 2s + 1$

$$\begin{aligned} u + jv &= F(s) = 2s + 1 = 2(\sigma + j\omega) + 1 \\ &\quad \swarrow \quad \uparrow \\ &\quad \text{scaling} \quad \text{shifting} \end{aligned}$$



$$\begin{aligned} &\text{s-plane} \\ &\sigma + j\omega = re^{j\varphi} \end{aligned}$$



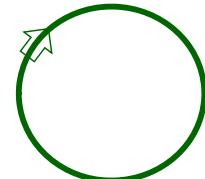
$$N = Z - P = 1 - 0 = 1$$

## Mapping Contours in the s-plane

- $F(s) : s \rightarrow F(s)$ , a conformal map, which retains the angles of the s-plane contour on the F(s) plane

- By convention, the area within a contour to the right of the traversal of the contour is considered to be the area enclosed by the contour

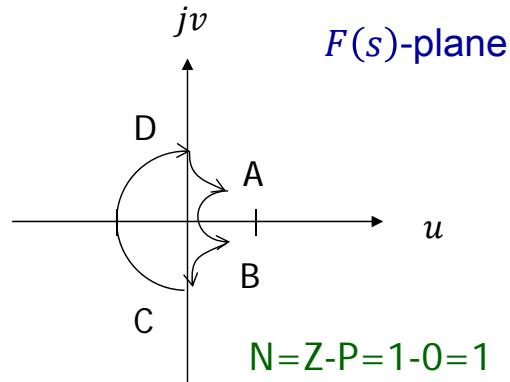
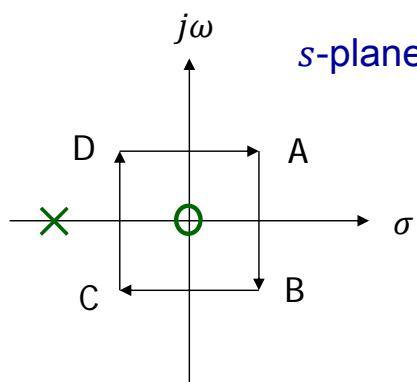
Clockwise : “+”



## Mapping Contours in the s-plane

- Example:  $F(s) = \frac{s}{s + 2}$

$$\begin{array}{ccccccc}
 A & \longrightarrow & B & \longrightarrow & C & \longrightarrow & D \\
 1+j1 & & 1 & & -j1 & & -1-j1 \\
 \frac{4+2j}{10} & & \frac{1}{3} & & \frac{4-2j}{10} & & \frac{1-2j}{5} \\
 & & & & & -j & -1 \\
 & & & & & & +j \\
 & & & & & & \frac{1+2j}{5}
 \end{array}$$



## Cauchy's Theorem -1

- Suppose :  $F(s) = \text{c.l. characteristic equ.}$

question : How to judge the stability of the closed-loop system given the open-loop transfer function

$L(s)$  ?

$$\Delta(s) = F(s) = 1 + L(s)$$

answer : Cauchy's Theorem & Nyquist Criterion

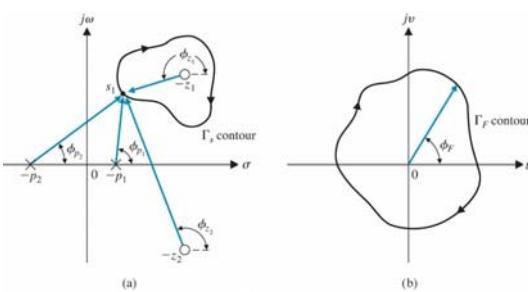
Note : in Chapter 7 課本符號不統一

$$\Delta(s) = 1 + KG(s) = 1 + F(s)$$

## Cauchy's Theorem -2

- If a contour  $\Gamma_s$  in the  $s$ -plane

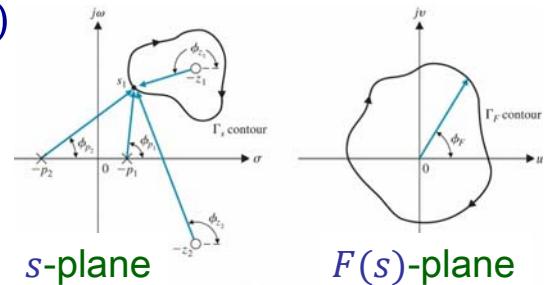
(1) encircles  $Z$  zeros and  $P$  poles of  $F(s)$ ,  
(2) does not pass through any poles or zeros of  $F(s)$ , and  
(3) the traversal is in the clockwise direction along the contour,  
the corresponding contour  $\Gamma_F$  in the  $F(s)$ -plane encircles the origin of the  $F(s)$ -plane  $N = Z - P$  times in the clockwise direction



## Cauchy's Theorem -3

□ Example:  $F(s) = \frac{(s + z_1)(s + z_2)}{(s + p_1)(s + p_2)}$

$$\begin{aligned} F(s) &= |F(s)| \angle F(s) \\ &= \frac{|s + z_1||s + z_2|}{|s + p_1||s + p_2|} (\angle(s + z_1) + \angle(s + z_2) - \angle(s + p_1) - \angle(s + p_2)) \\ &= |F(s)|(\angle\phi_{z_1} + \angle\phi_{z_2} - \angle\phi_{p_1} - \angle\phi_{p_2}) \end{aligned}$$



As  $s$  traverses the contour  $\Gamma_s$  (a full rotation)

$\phi_{z_2}, \phi_{p_1}, \phi_{p_2}$ : net angle change = 0

$\phi_{z_1}$ : net angle change =  $360^\circ$

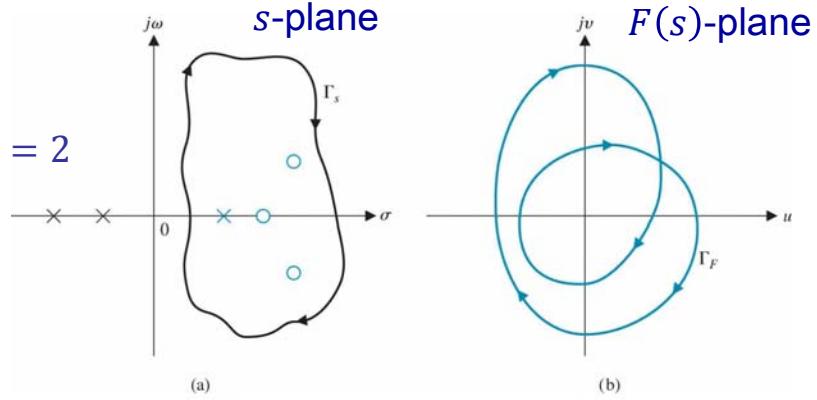
If  $F(s)$  has  $Z$  zeros and  $P$  poles,  $\phi_z = 2\pi(Z)$   $\phi_p = 2\pi(P)$

$$\begin{aligned} \phi_F &= \phi_Z - \phi_P \\ \Rightarrow N &= Z - P \end{aligned}$$

## Cauchy's Theorem -4

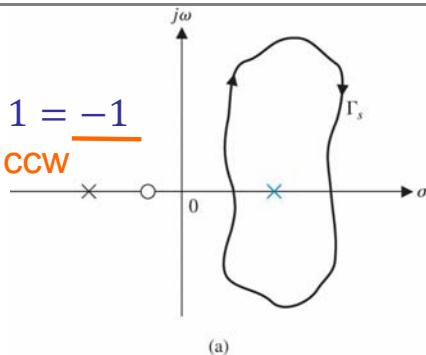
□ Examples:

$$N = Z - P = 3 - 1 = 2$$



$$N = Z - P = 0 - 1 = -1$$

CCW



# Nyquist Criterion -1

## □ Logic

$$L(s) = \frac{N(s)}{D(s)}: \text{Loop T.F., known}$$

$$T(s) = \frac{\dots}{\Delta(s)}: \text{Closed-loop T.F.}$$

$$\Delta(s) = F(s) = \frac{k \prod_{i=1}^N (s + s_i)}{\prod_{k=1}^M (s + s_k)} = 1 + L(s) = 1 + \frac{N(s)}{D(s)} = \frac{D(s) + N(s)}{D(s)}$$

poles of  $F(s)$  = poles of  $L(s)$  = roots of  $D(s)$  known

zeros of  $F(s)$  = poles of  $T(s)$  = roots of  $D(s) + N(s)$  unknown

Determining stability of the system

# Nyquist Criterion -2

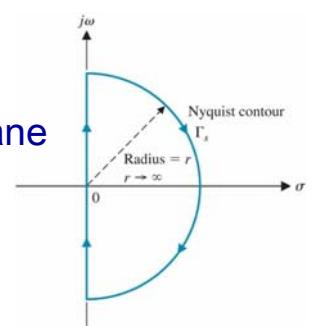
A stable c.l. system  $\Rightarrow$  all “poles” of  $T(s)$  in LHP  
 $\Rightarrow$  all “zeros” of  $F(s)$  in LHP



Choose a contour  $\Gamma_s$  encloses the entire RHP in  $s$ -plane



Cauchy's Theorem       $Z = N + P$



knowing  $N$  by  $\Gamma_F$ 's encirclement of  $(0,0)$  in  $F(s)$ -plane  
 or  $\Gamma_L$ 's encirclement of  $(-1,0)$  in  $L(s)$ -plane

knowing  $P$  by  $L(s)$  its poles

$\Rightarrow$  Know  $Z$  zeros of  $F(s)$  in RHP

$\Rightarrow$  Know  $Z$  poles of  $T(s)$  in RHP

## Nyquist Criterion -3

- A feedback system is stable if and only if the contour  $\Gamma_L$  in the  $L(s)$ -plane does NOT encircle the  $(-1,0)$  point when the number of poles of  $L(s)$  in the right-hand s-plane is zero ( $P = 0$ )

$$Z = N + P = 0 + 0 = 0$$

- A feedback control system is stable if and only if, for the contour  $\Gamma_L$ , the number of counterclockwise encirclement of the  $(-1,0)$  point is equal to the number of poles of  $L(s)$  with positive real parts

$$Z = N + P = 0$$

### Example 1 -1

- $L(s) = GH(s) = \frac{K}{(\tau_1 s + 1)(\tau_2 s + 1)} \quad K > 0$

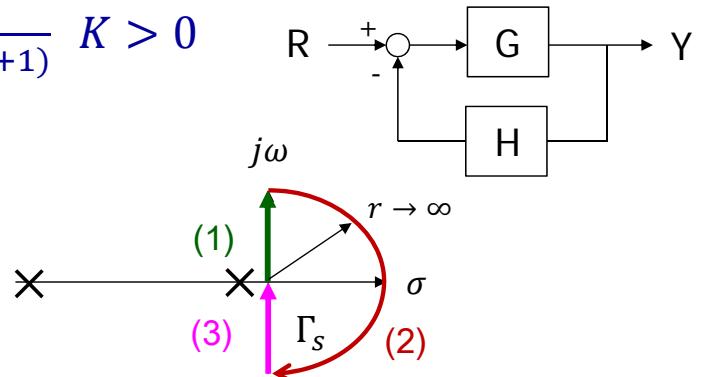
Assume  $\tau_1 = 1 \quad \tau_2 = \frac{1}{10}$

(1)  $\omega = 0 \rightarrow \omega = +\infty$

$$GH(j\omega) = GH(s) \Big|_{s=j\omega}$$

$$= \frac{10K(10 - \omega^2)}{(10 - \omega^2)^2 - (11\omega^2)^2} + j \frac{-10K(11\omega)}{(10 - \omega^2)^2 - (11\omega^2)^2}$$

Cross Im-axis at  $|Im|_{\omega=\sqrt{10}} = 0.287K$  when  $\omega = \sqrt{10}$



$$|GH| = 10K \sqrt{\frac{1}{(10 - \omega^2)^2 + (11\omega^2)^2}}$$

$$\phi = -\tan^{-1}\left(\frac{11\omega}{10 - \omega^2}\right)$$

"4 quadrants"

## Example 1 -2

$\omega$	0	1	$\sqrt{10}$	10	100	$\infty$
$ GH $	100	70.7	28.74	6.8	0.1	0
$\phi$	0	-50.7	-90	-129.3	-173.7	-180

P.S. The “polar plot” in Chapter 8

$$(2) \omega = +\infty \rightarrow \omega = -\infty$$

$$s = re^{j\phi} : \begin{cases} \phi = 90^\circ \rightarrow -90^\circ \text{ cw} \\ r \rightarrow \infty \end{cases}$$

$$L(s) = le^{j\theta} = \lim_{r \rightarrow \infty} GH(s) \Big|_{s=re^{j\phi}} = \lim_{r \rightarrow \infty} \left| \frac{K}{\tau_1 \tau_2 r^2} \right| e^{-j2\phi}$$

$$\begin{cases} \theta = -180^\circ \rightarrow +180^\circ \text{ ccw} \\ l \rightarrow 0 \end{cases} \quad \curvearrowleft$$

P. S.  $\rightarrow 0$  when the denominator has higher order than the numerator

## Example 1 -3

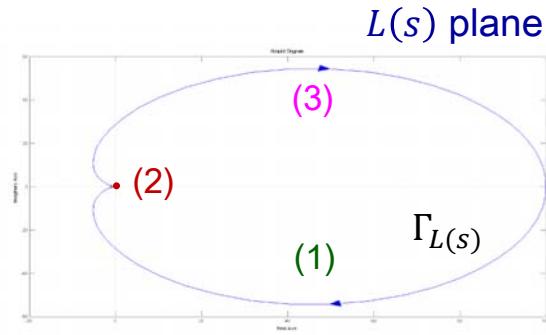
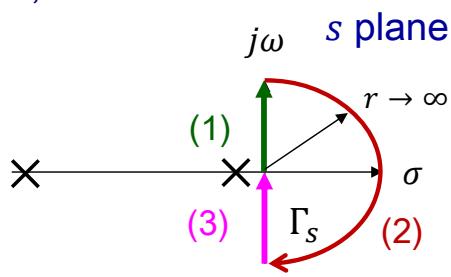
$$(3) \omega = -\infty \rightarrow \omega = 0$$

$$L(s): GH(s) \Big|_{s=-j\omega} = GH(-j\omega) = \text{complex conjugate of } GH(j\omega)$$

$$\omega = +\infty \rightarrow \omega = 0^+$$

→ mirrored of (1) w.r.t. Re-axis  
and change arrow direction

Thus,



$$Z = N + P = 0 + 0 = 0 \text{ stable}, \quad K \text{ doesn't matter}$$

## Example 2 -1

□  $L(s) = GH(s) = \frac{K}{s(\tau s + 1)}$   $K > 0, \tau > 0$

pole at origin  $\rightarrow \Gamma_s$  needs detour

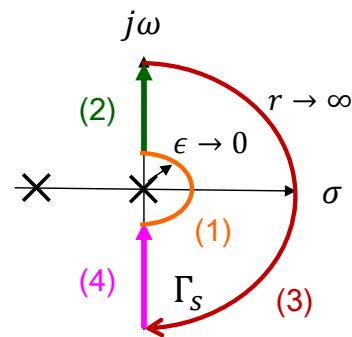
(1)  $\omega = 0^- \rightarrow \omega = 0^+$

$$s = re^{j\phi} : \begin{cases} \phi = -90^\circ \rightarrow 90^\circ \text{ ccw} \\ r = \epsilon \rightarrow 0 \end{cases}$$

$$L(s) = le^{j\theta} = \lim_{\epsilon \rightarrow 0} GH(s) \approx \lim_{\epsilon \rightarrow 0} \left( \frac{K}{\epsilon e^{j\phi}} \right) = \lim_{\epsilon \rightarrow 0} \left( \frac{K}{\epsilon} \right) e^{-j\phi}$$

$$\begin{cases} \theta = 90^\circ \rightarrow -90^\circ \text{ cw} \\ l \rightarrow \infty \end{cases}$$

P.S. an infinite half circle



## Example 2 -2

(2)  $\omega = 0^+ \rightarrow \omega = +\infty$

The same as the “polar plot” example shown in Chap. 8

$$GH(j\omega) = \frac{K}{j\omega(j\omega\tau + 1)} = \frac{K}{-\omega^2\tau + j\omega} = \frac{-K\omega^2\tau}{\omega^4\tau^2 + \omega^2} + \frac{-j\omega K}{\omega^4\tau^2 + \omega^2}$$

$\pm\omega \rightarrow \infty$ : symmetry w.r.t. Re-axis

	<b>R</b>	<b>X</b>		<b>G</b>	<b>ϕ</b>
$\omega = 0$	$-K\tau$	$-\infty$		$\omega = 0$	$\infty$
$\omega = \frac{1}{2}$	$-\frac{K\tau}{2}$	$-\frac{K\tau}{2}$		$\omega = \frac{1}{2}$	$-135^\circ$
$\omega = \infty$	0	0		$\omega = \infty$	$180^\circ$

$$|GH| = \frac{K}{(\omega^4\tau^2 + \omega^2)^{\frac{1}{2}}} \quad \text{“4 quadrants”}$$

$$\phi(\omega) = -\tan^{-1} \left( \frac{1}{-\omega\tau} \right) \text{ or } = -\frac{\pi}{2} - \tan^{-1} \omega\tau$$

## Example 2 -3

(3)  $\omega = +\infty \rightarrow \omega = -\infty$

$$s = re^{j\phi} : \begin{cases} \phi = 90^\circ \rightarrow -90^\circ \text{ cw} \\ r \rightarrow \infty \end{cases}$$

$$L(s) = le^{j\theta} = \lim_{r \rightarrow \infty} GH(s) \Big|_{s=re^{j\phi}} = \lim_{r \rightarrow \infty} \left| \frac{K}{\tau r^2} \right| e^{-j2\phi}$$

$$\begin{cases} \theta = -180^\circ \rightarrow +180^\circ \text{ ccw} \\ l \rightarrow 0 \end{cases} \quad \curvearrowleft$$

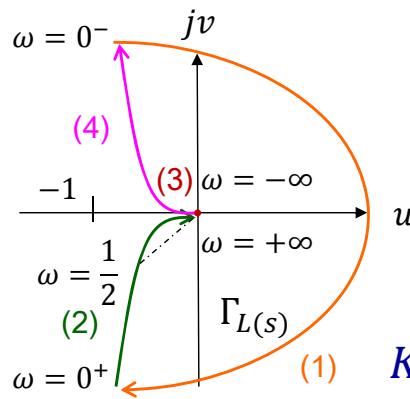
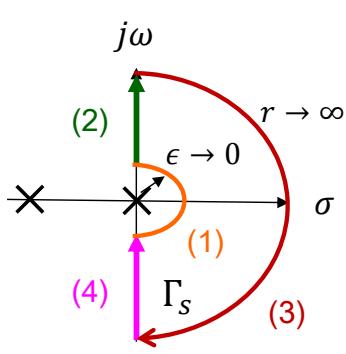
(4)  $\omega = -\infty \rightarrow \omega = 0^-$

$L(s)$ :  $GH(s) \Big|_{s=-j\omega} = GH(-j\omega) = \text{complex conjugate of } GH(j\omega)$   
 $\omega = +\infty \rightarrow \omega = 0^+$

→ mirrored of (1) w.r.t. Re-axis  
and change arrow direction

## Example 2 -4

Thus,



$Z = N + P = 0 + 0 = 0$   
→ stable  
 $K$  doesn't matter

Note: Routh-Hurwitz & Root locus methods

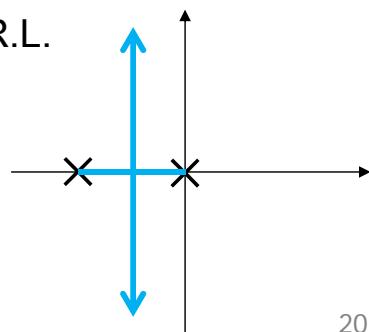
$$L(s) = GH(s) = \frac{K}{s(\tau s + 1)} \quad K > 0$$

R.H.

$$\tau s^2 + s + K = 0$$

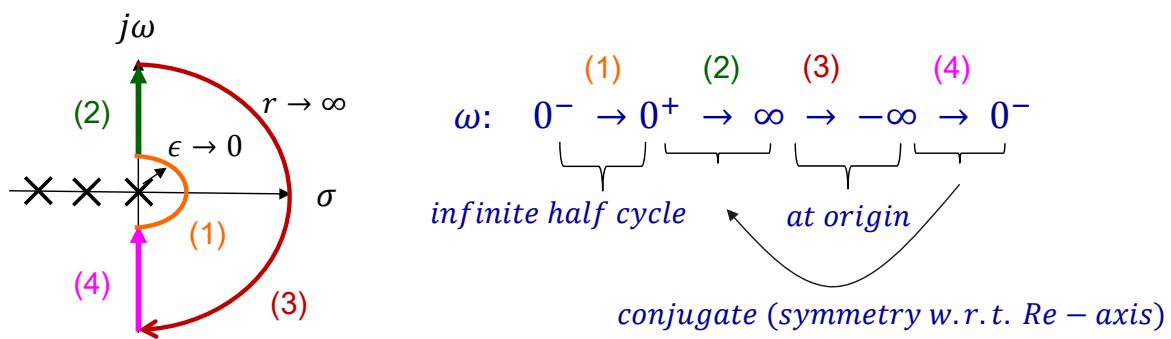
Stable as long as  $K > 0$

R.L.



## Example 3 -1

□  $GH(s) = \frac{K}{s(\tau_1 s + 1)(\tau_2 s + 1)} \quad K > 0, \tau_1 > 0, \tau_2 > 0$



(1)  $\omega = 0^- \rightarrow \omega = 0^+$

$$s = re^{j\phi} : \begin{cases} \phi = -90^\circ \rightarrow 90^\circ \text{ ccw} \\ r = \epsilon \rightarrow 0 \end{cases}$$

$$L(s) = le^{j\theta} = \lim_{\epsilon \rightarrow 0} \frac{K}{\epsilon e^{j\phi} (\tau_1 \epsilon e^{j\phi} + 1)(\tau_2 \epsilon e^{j\phi} + 1)} \approx \lim_{\epsilon \rightarrow 0} \left( \frac{K}{\epsilon e^{j\phi}} \right)$$

$$\begin{cases} \theta = 90^\circ \rightarrow -90^\circ \text{ cw} \\ l \rightarrow \infty \end{cases} \quad \text{P.S. an infinite half circle}$$

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## Example 3 -2

(2)  $\omega = 0^+ \rightarrow \omega = +\infty$

$$GH(j\omega) = \frac{-K(\tau_1 + \tau_2) - jK\left(\frac{1}{\omega}\right)(1 - \omega^2\tau_1\tau_2)}{1 + \omega^2(\tau_1^2 + \tau_2^2) + \omega^4\tau_1^2\tau_2^2}$$

$$\text{when } \omega^2 = \frac{1}{\tau_1\tau_2}, \quad X = v = \frac{K \frac{1}{\omega} (1 - \omega^2\tau_1\tau_2)}{\sim} = 0$$

across  $Re - axis$  at

$$R = u = \left. \frac{-K(\tau_1 + \tau_2)}{1 + \omega^2(\tau_1^2 + \tau_2^2) + \omega^4\tau_1\tau_2} \right|_{\omega^2 = \frac{1}{\tau_1\tau_2}} = \frac{-K\tau_1\tau_2}{\tau_1 + \tau_2}$$

## Example 3 -3

(3)  $\omega = +\infty \rightarrow \omega = -\infty$

$$s = re^{j\phi} : \begin{cases} \phi = 90^\circ \rightarrow -90^\circ \text{ cw} \\ r \rightarrow \infty \end{cases}$$

$$L(s) = le^{j\theta} = \lim_{r \rightarrow \infty} \frac{K}{re^{j\phi}(\tau_1 re^{j\phi} + 1)(\tau_2 re^{j\phi} + 1)} = \Bigg|_{s=re^{j\phi}}$$

$\underbrace{\sim \tau_1 re^{j\phi}}_{\sim \tau_1 \tau_2 r^3} \quad \underbrace{\sim \tau_2 re^{j\phi}}$

$$\approx \lim_{r \rightarrow \infty} \left| \frac{K}{\tau_1 \tau_2 r^3} \right| e^{-j3\phi}$$

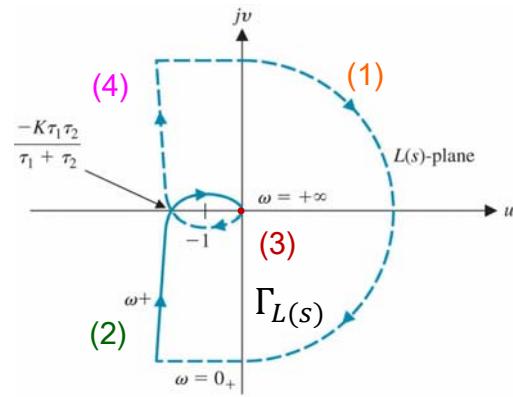
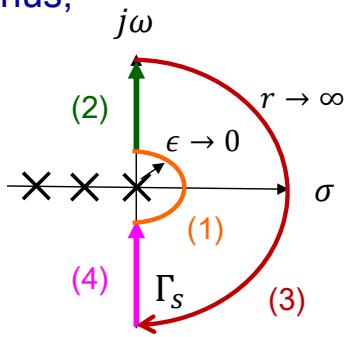
$$\begin{cases} \theta = -270^\circ \rightarrow +270^\circ \text{ ccw} \\ l \rightarrow 0 \end{cases} \quad \curvearrowleft$$

(4)  $\omega = -\infty \rightarrow \omega = 0^-$

*mirrored of (1) w.r.t. Re-axis and change arrow direction*

## Example 3 -4

Thus,



when  $\frac{-K\tau_1\tau_2}{\tau_1+\tau_2} > -1, N = 0, Z = N + P = 0 + 0 = 0 \text{ stable}$

$$K < \frac{\tau_1 + \tau_2}{\tau_1 \tau_2}$$

when  $\frac{-K\tau_1\tau_2}{\tau_1+\tau_2} < -1, N = 2, Z = 2 + 0 = 2 \text{ unstable}$

"2 RHP poles"

## Example 3 -5

Note: Routh-Hurwitz & Root locus methods

$$GH(s) = \frac{K}{s(\tau_1 s + 1)(\tau_2 s + 1)}$$

R.H.

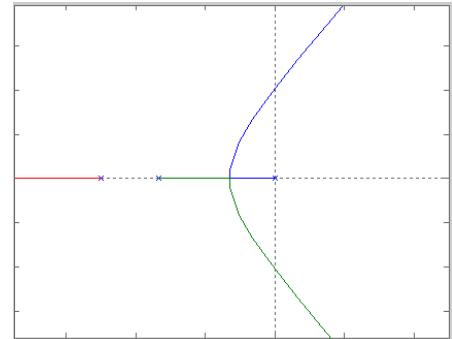
$$\tau_1 \tau_2 s^3 + (\tau_1 + \tau_2) s^2 + s + K = 0$$

$$\textcircled{1} \quad K > 0$$

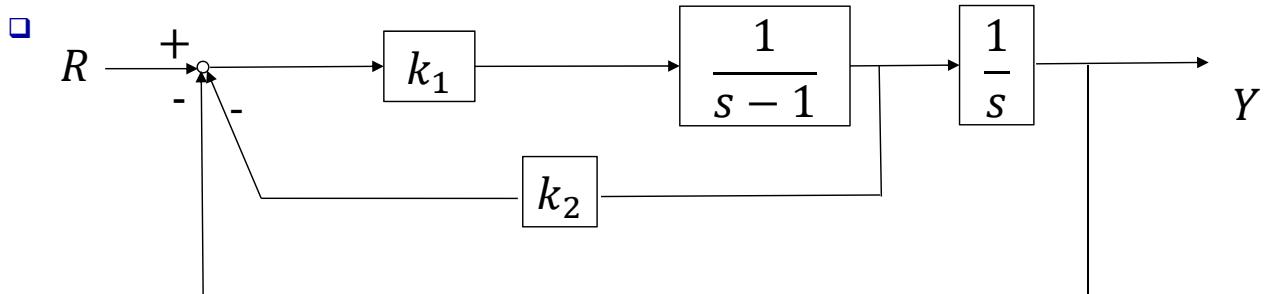
$$\textcircled{2} \quad (\tau_1 + \tau_2) > \tau_1 \tau_2 K$$

$$\rightarrow K < \frac{\tau_1 + \tau_2}{\tau_1 \tau_2}$$

R.L.



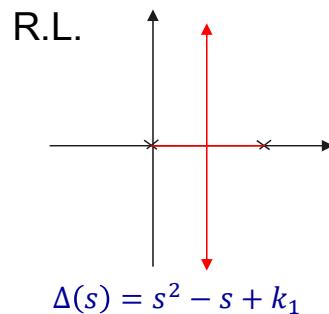
## Example 4 -1



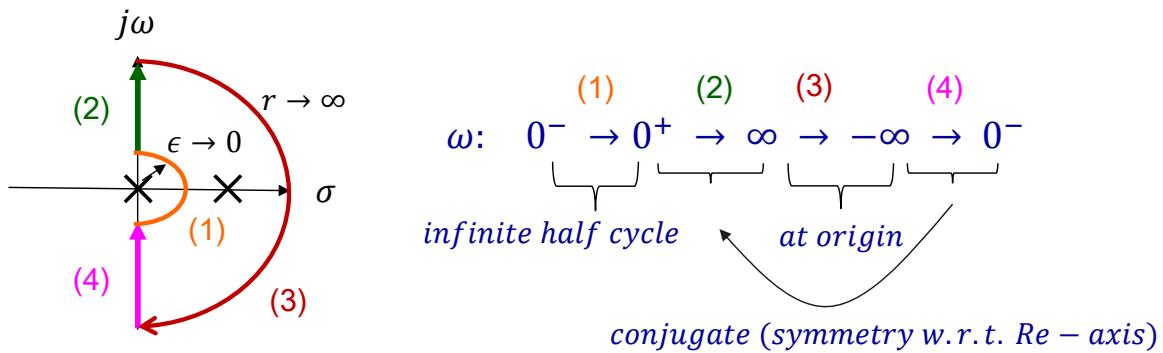
$k_2 = 0$ , without derivative feedback

$$GH(s) = \frac{k_1}{s(s-1)}$$

$P = 1$ ,  
stable system ( $Z = 0$ ):  $N = -P = -1$



## Example 4 -2



$$(1) \omega = 0^- \rightarrow \omega = 0^+$$

$$s = re^{j\phi} : \begin{cases} \phi = -90^\circ \rightarrow 90^\circ & ccw \\ r = \epsilon \rightarrow 0 & \end{cases}$$

$$L(s) = le^{j\theta} = \lim_{\epsilon \rightarrow 0} \frac{k_1}{\epsilon e^{j\phi}(\epsilon e^{j\phi} - 1)} \approx \lim_{\epsilon \rightarrow 0} \left( \frac{k_1}{-\epsilon e^{j\phi}} \right) = \lim_{\epsilon \rightarrow 0} \left( \frac{k_1}{\epsilon} \right) e^{j(-180-\phi)}$$

$$\begin{cases} \theta = -90^\circ \rightarrow -270^\circ & cw \\ l \rightarrow \infty & \end{cases}$$

P.S. an infinite half circle

## Example 4 -3

$$(2) \omega = 0^+ \rightarrow \omega = +\infty$$

$$GH(j\omega) = \frac{k_1}{j\omega(j\omega - 1)} = \frac{-k_1\omega^2\tau + jk_1\omega}{\omega^4\tau^2 + \omega^2}$$

$$(3) \omega = +\infty \rightarrow \omega = -\infty$$

$$s = re^{j\phi} : \begin{cases} \phi = 90^\circ \rightarrow -90^\circ & cw \\ r \rightarrow \infty & \end{cases}$$

$$L(s) = le^{j\theta} = \left| \underbrace{\lim_{r \rightarrow \infty} \frac{k_1}{re^{j\phi}(re^{j\phi} - 1)}}_{\sim re^{j\phi}} \right|_{s=re^{j\phi}} \approx \lim_{r \rightarrow \infty} \left| \frac{K}{r^2} \right| e^{-j2\phi}$$

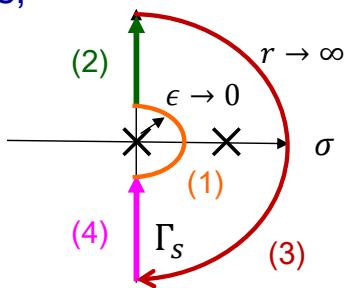
$$\begin{cases} \theta = -180^\circ \rightarrow +180^\circ & ccw \\ l \rightarrow 0 & \end{cases}$$

$$(4) \omega = -\infty \rightarrow \omega = 0^-$$

mirrored of (1) w.r.t. Re-axis and change arrow direction

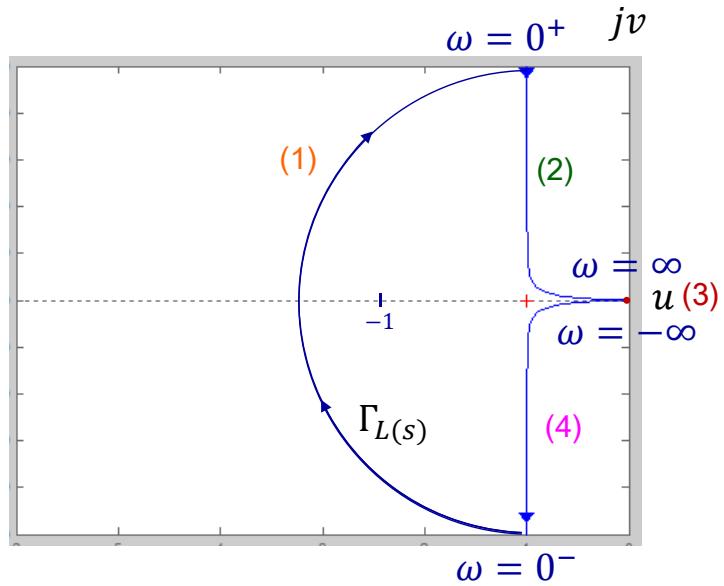
## Example 4 - 4

Thus,



$$\frac{|GM|}{k_1} \quad \angle GM$$

$j0^-$	$\infty$	$-90^\circ$
$j0^+$	$\infty$	$90^\circ$
$j$	$\frac{1}{\sqrt{2}}$	$135^\circ$
$+j\infty$	0	$180^\circ$
$-j\infty$	0	$-180^\circ$

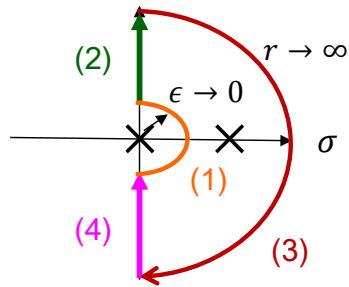


$\therefore N = 1 \ Z = 2, \text{ unstable}$

## Example 4 - 5

□ with  $k_2$

$$GH(s) = \frac{k_1(1 + k_2 s)}{s(s-1)}$$



$$(1) \omega = 0^- \rightarrow \omega = 0^+$$

$$s = re^{j\phi} : \begin{cases} \phi = -90^\circ \rightarrow 90^\circ & \text{ccw} \\ r = \epsilon \rightarrow 0 & \end{cases}$$

$$L(s) = le^{j\theta} = \lim_{\epsilon \rightarrow 0} \frac{k_1(1 + k_2 \epsilon e^{j\phi})}{\epsilon e^{j\phi}(\epsilon e^{j\phi} - 1)} \approx \lim_{\epsilon \rightarrow 0} \left( \frac{k_1}{-\epsilon e^{j\phi}} \right) = \lim_{\epsilon \rightarrow 0} \left( \frac{k_1}{\epsilon} \right) e^{j(-180 - \phi)}$$

$$\begin{cases} \theta = -90^\circ \rightarrow -270^\circ & \text{cw} \\ l \rightarrow \infty & \end{cases}$$

## Example 4 -6

(2)  $\omega = 0^+ \rightarrow \omega = +\infty$

$$GH(j\omega) = \frac{k_1(1 + k_2 j\omega)}{-\omega^2 - j\omega} = -\frac{k_1(\omega^2 + \omega^2 k_2)}{\omega^2 + \omega^4} + \frac{j(\omega - k_2 \omega^3)k_1}{\omega^2 + \omega^4}$$

$\omega - k_2 \omega^3 = 0$  across Re-axis

$$u|_{\omega^2 = \frac{1}{k_2}} = -\frac{k_1(\omega^2 + \omega^2 k_2)}{\omega^2 + \omega^4}|_{\omega^2 = \frac{1}{k_2}} = -k_1 k_2$$

(3)  $\omega = +\infty \rightarrow \omega = -\infty$

$$s = re^{j\phi} : \begin{cases} \phi = 90^\circ \rightarrow -90^\circ \text{ cw} \\ r \rightarrow \infty \end{cases}$$

$$L(s) = le^{j\theta} = \left| \lim_{r \rightarrow \infty} \frac{k_1(1 + k_2 re^{j\phi})}{re^{j\phi}(re^{j\phi} - 1)} \right|_{s=re^{j\phi}} \approx \lim_{r \rightarrow \infty} \left| \frac{k_1 k_2}{r} \right| e^{-j\phi}$$

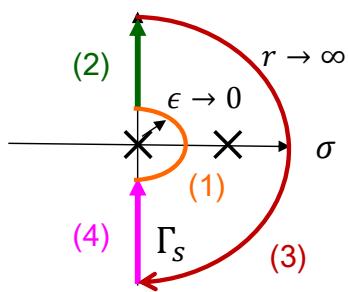
$$\begin{cases} \theta = -90^\circ \rightarrow +90^\circ \text{ ccw} \\ l \rightarrow 0 \end{cases}$$

## Example 4 -7

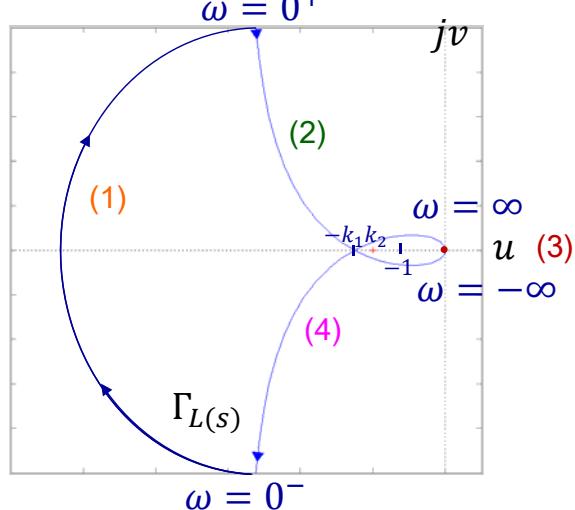
(4)  $\omega = -\infty \rightarrow \omega = 0^-$

mirrored of (1) w.r.t. Re-axis and change arrow direction

Thus,



when  $k_1 k_2 > 1$



$\Gamma_{F(s)}$  encircles  $-1$  once in ccw direction

$N = -1, Z = N + P = -1 + 1 = 0,$   
stable

$$\begin{aligned} \Delta(s) &= s^2 - s + k_1 k_2 s + k_1 \\ &= s^2 + (k_1 k_2 - 1)s + k_1 = 0 \\ k_1 k_2 - 1 > 0 &\quad k_1 > 0 \end{aligned}$$

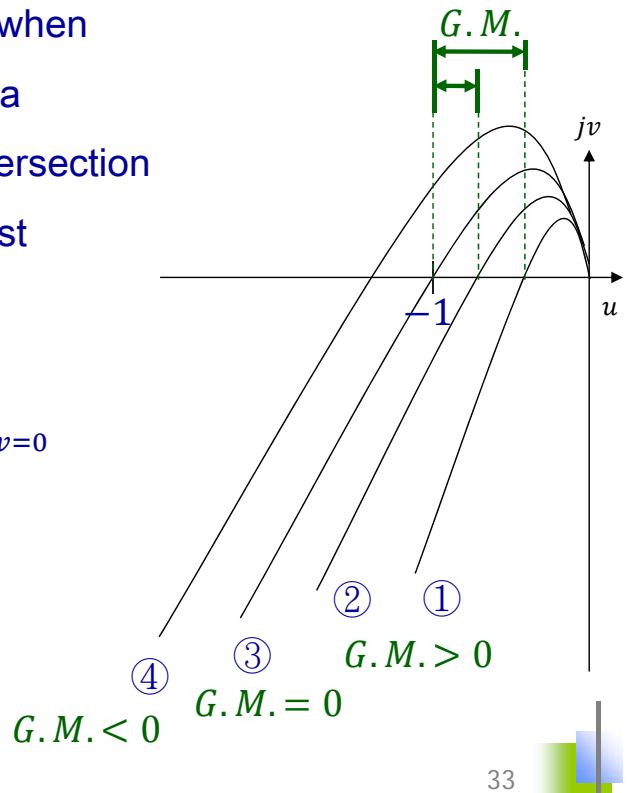
# Gain Margin and Phase Margin -1

## □ Gain margin

- The increase in the system gain when phase =  $-180^\circ$  that will result in a marginally stable system with intersection of the  $-1 + j0$  point on the Nyquist diagram

$$\begin{aligned} G.M. &\triangleq 20 \log |1| - 20 \log |L(\omega)|_{v=0} \\ &= 20 \log \frac{1}{|L(\omega)|_{v=0}} \text{ dB} \end{aligned}$$

$$G.M. = 0 - G_{dB}|_{v=0} \text{ dB}$$

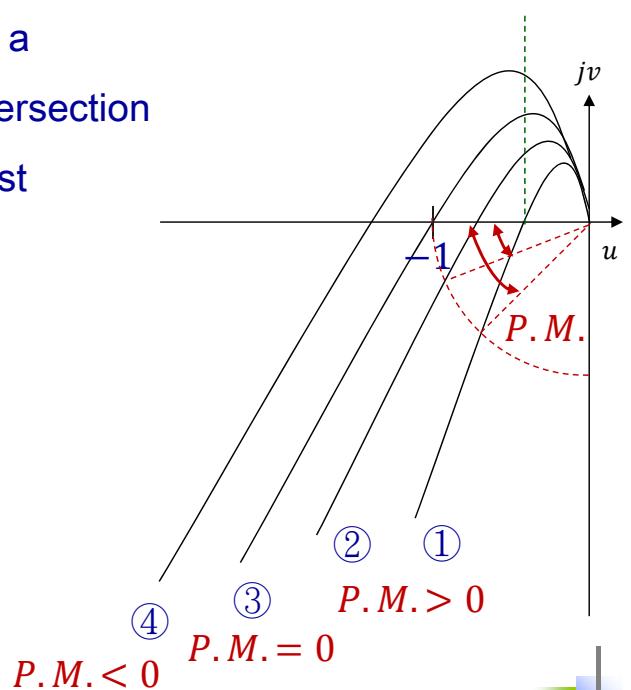


# Gain Margin and Phase Margin -2

## □ Phase margin

- The amount of phase shift of the  $L(j\omega)$  at unity magnitude that will result in a marginally stable system with intersection of the  $-1 + j0$  point on the Nyquist diagram

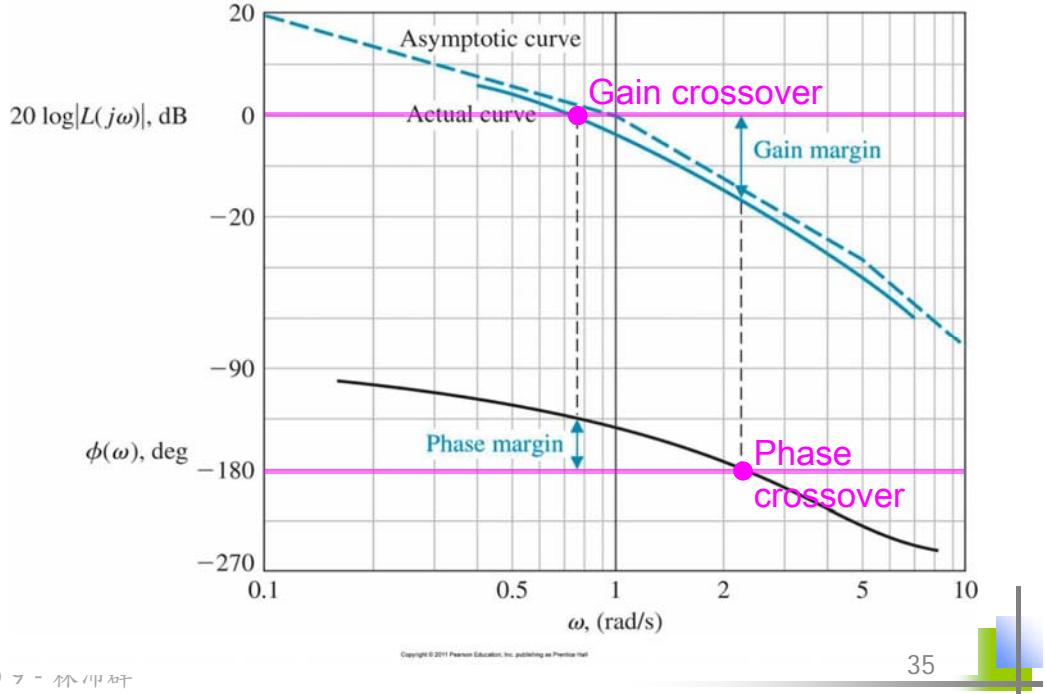
$$P.M. = \phi_{PM} = \angle L(\omega) - (-180^\circ)$$



## Gain Margin and Phase Margin -3

### □ Bode plot

- ◆ Gain / phase crossover & gain / phase margin are more easily determined (on the Bode plot than on the Nyquist plot)

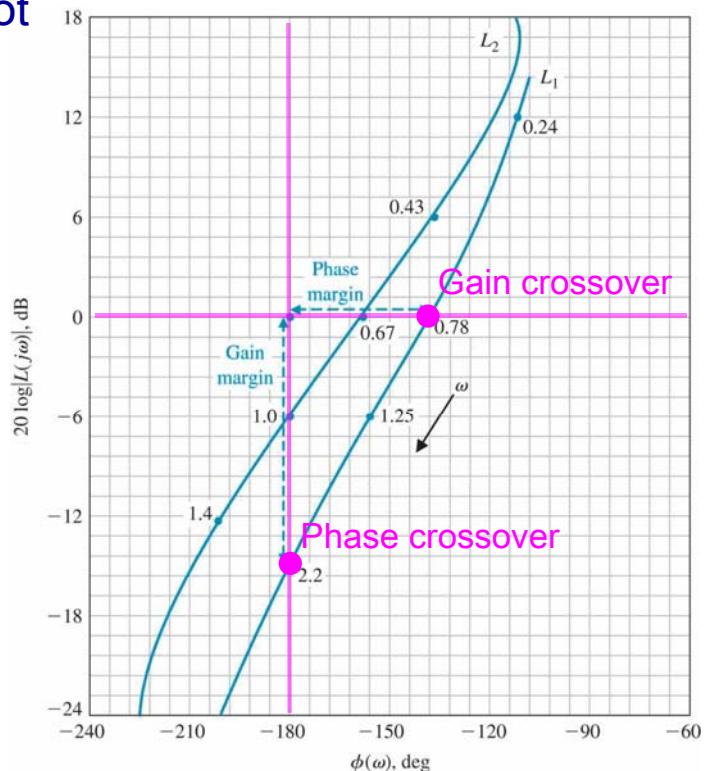


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## Gain Margin and Phase Margin -4

### □ Log-magnitude-phase plot



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## Gain Margin and Phase Margin -5

- Ex: A standard 2<sup>nd</sup>-order system

$$L(s) = GH(s) = \frac{\omega_n^2}{s(s + 2\xi\omega_n)}$$

$$GH(j\omega) = \frac{\omega_n^2}{j\omega(j\omega + 2\xi\omega_n)}$$

$$\left|GH(\omega_g)\right| = 1 = \frac{\omega_n^2}{\omega_g(\omega_g^2 + 4\xi^2\omega_n^2)^{\frac{1}{2}}}$$

↗ Gain crossover

$$(\omega_g^2)^2 + 4\xi^2\omega_n^2(\omega_g^2) - \omega_n^4 = 0$$

$$\left(\frac{\omega_g^2}{\omega_n^2}\right)^2 + 4\xi^2\left(\frac{\omega_g^2}{\omega_n^2}\right) - 1 = 0$$

$$\Rightarrow \frac{\omega_g^2}{\omega_n^2} = (4\xi^4 + 1)^{\frac{1}{2}} - 2\xi^2$$

## Gain Margin and Phase Margin -6

$$P.M. = \phi_{PM} = -90^\circ - \tan^{-1}\left(\frac{\omega_g}{2\xi\omega_n}\right) - (-180^\circ)$$

$$= \tan^{-1}(2\xi \left[ \frac{1}{(4\xi^4 + 1)^{\frac{1}{2}} - 2\xi^2} \right]^{\frac{1}{2}})$$

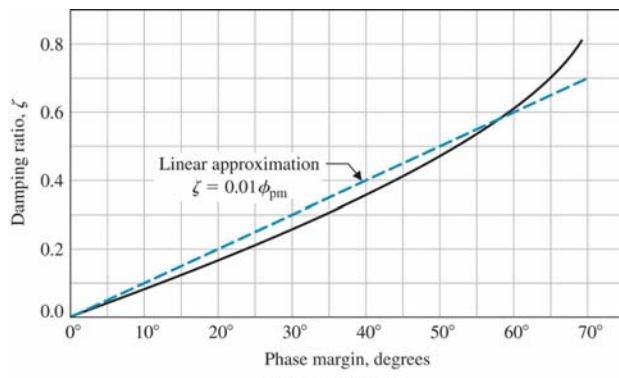
$$= \tan^{-1}\left(\frac{2}{\left[\left(4 + \frac{1}{\xi^2}\right)^{\frac{1}{2}} - 2\right]^{\frac{1}{2}}}\right)$$

approximation:

$$\xi = 0.01\phi_{PM} \quad \xi \leq 0.7$$

Adjust  $\phi_{PM}$  in frequency response is EQUAL to adjust  $\xi$  in time response

$G.M. \rightarrow \infty$



## The O.I. T.F. vs. C.I. T.F. -1

- Question: Can we obtain closed-loop frequency response from the open-loop frequency response?

Assuming unity feedback  $H(j\omega) = 1$

$$G_c G(j\omega) = u + jv$$

$$\text{Closed-loop T.F. } T(j\omega) = \frac{G_c G(j\omega)}{1+G_c G(j\omega)} = \frac{u+jv}{(1+u)+jv} = M(\omega)e^{j\phi(\omega)}$$

$$M(\omega) = \left| \frac{G_c G(j\omega)}{1+G_c G(j\omega)} \right| = \left| \frac{u+jv}{1+u+jv} \right| = \frac{(u^2+v^2)^{\frac{1}{2}}}{[(1+u)^2+v^2]^{\frac{1}{2}}}$$

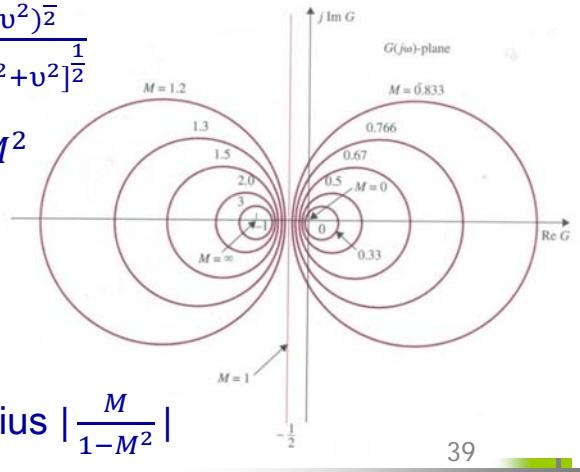
$$(1-M^2)u^2 + (1-M^2)v^2 - 2M^2u = M^2$$

$$u^2 + v^2 - \frac{2M^2}{1-M^2}u = \frac{M^2}{1-M^2}$$

→  $(u - \frac{M^2}{1-M^2})^2 + v^2 = (\frac{M}{1-M^2})^2$

A circle: center at  $(\frac{M^2}{1-M^2}, 0)$ , radius  $|\frac{M}{1-M^2}|$

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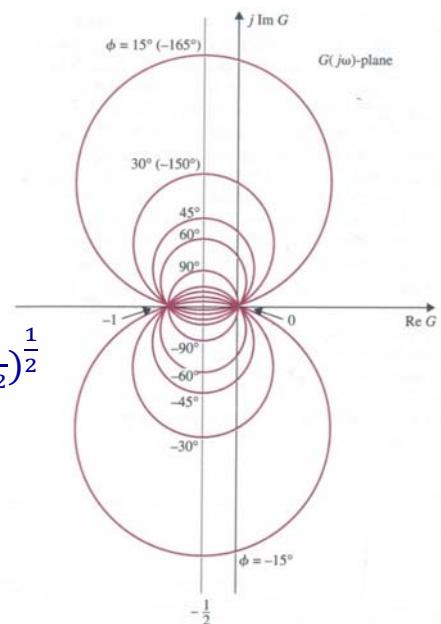
## The O.I. T.F. vs. C.I. T.F. -2

$$\tan(\phi(\omega)) = N = \frac{v}{u+u^2+v^2}$$

$$u^2 + v^2 + u - \frac{v}{N} = 0$$

$$\left(u + \frac{1}{2}\right)^2 + \left(v - \frac{1}{2N}\right)^2 = \frac{1}{4}\left(1 + \frac{1}{N^2}\right)$$

A circle: center at  $(-\frac{1}{2}, \frac{1}{2N})$ , radius  $\frac{1}{2}(1 + \frac{1}{N^2})^{\frac{1}{2}}$



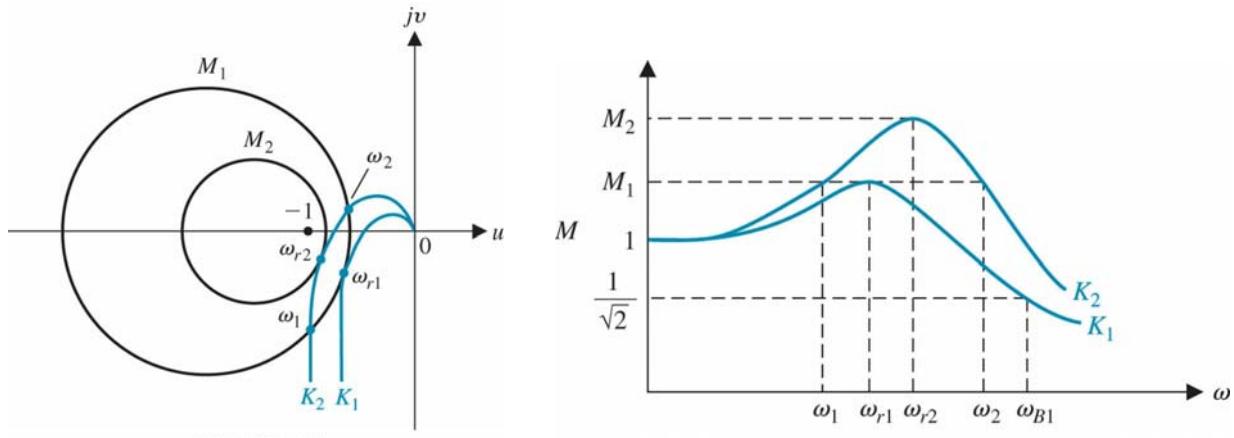
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# The O.I. T.F. vs. C.I. T.F. -3

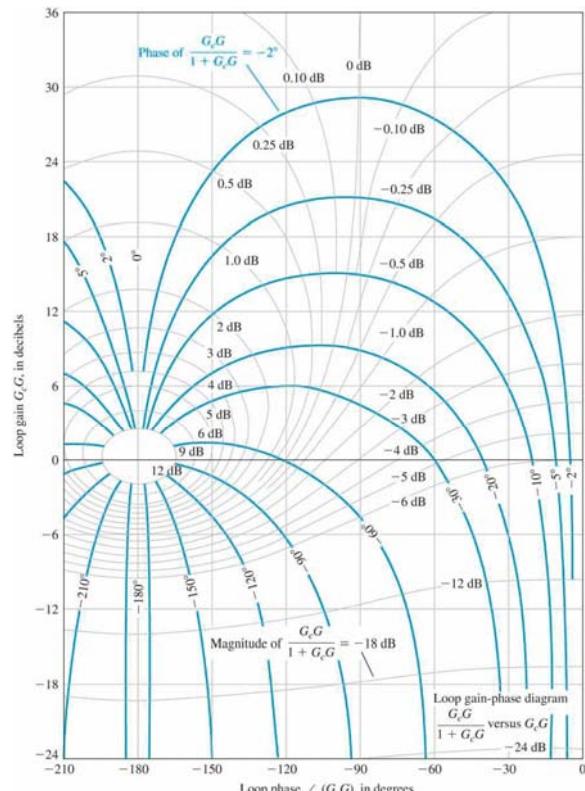
- Ex: A system with two different gains,  $K_1$  &  $K_2$

由open-loop T.F.的polar plot軌跡和M圓軌跡相交的狀態，可以推估出此系統在closed-loop後的frequency response狀態



## Nichols Chart -1

- Plotting magnitude and phase of the closed-loop system as contours on the log-magnitude-phase diagram



# Nichols Chart -2

□ Ex:  $G(j\omega) = \frac{0.64}{j\omega[(j\omega)^2 + j\omega + 1]}$   
 $H(j\omega) = 1$

$$\zeta = 0.5$$

$$P.M. = 30^\circ$$

C.I.  $20\log M_{p\omega} = 9 \text{ dB}$   
at  $\omega_r = 0.9$

